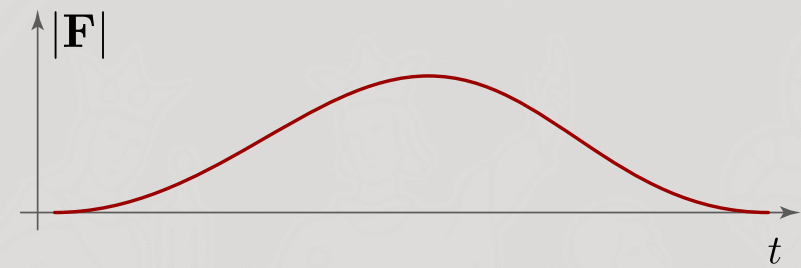
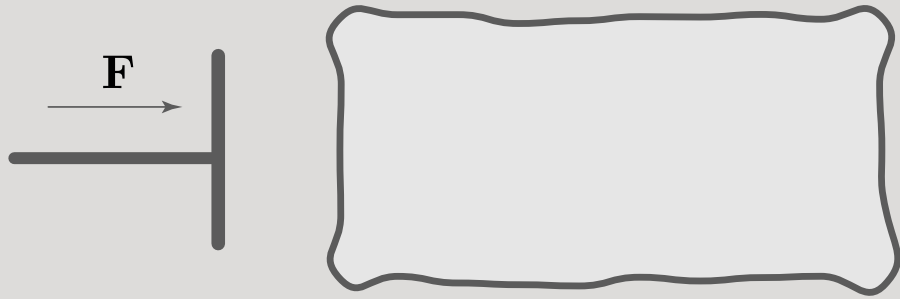
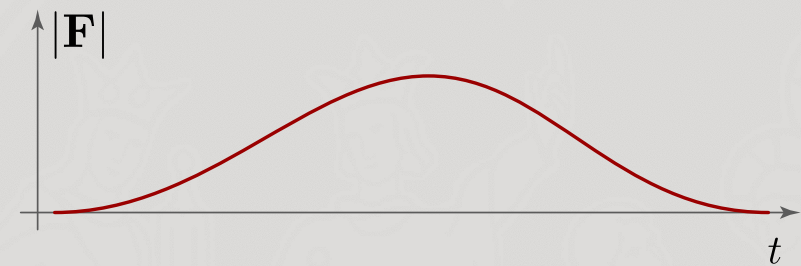
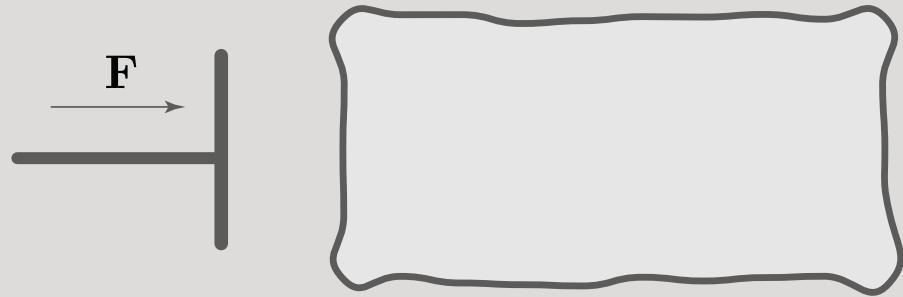


example: Jarzynski relation



Work done: W
Free energy gain: ΔF
thermodynamics: $\langle W \rangle \geq \Delta F$

example: Jarzynski relation

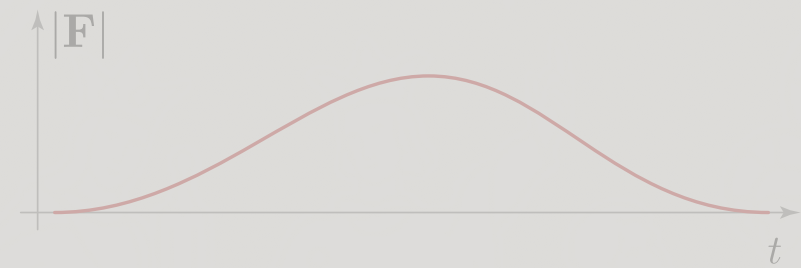


Work done: W
Free energy gain: ΔF
thermodynamics: $\langle W \rangle \geq \Delta F$

Jarzynski relation (1997):

$$\left\langle e^{-\frac{1}{T}(W - \Delta F)} \right\rangle = 1$$

implies existence of processes in which **work is gained out of fluctuations!**



FR: describe the interplay of fluctuations and

Work done: W
Free energy gain: ΔF
dissipation out of equilibrium

thermodynamics: $\langle W \rangle \geq \Delta F$

Jarzynski relation (1997):

$$\left\langle e^{-\frac{1}{T} (W - \Delta F)} \right\rangle = 1$$

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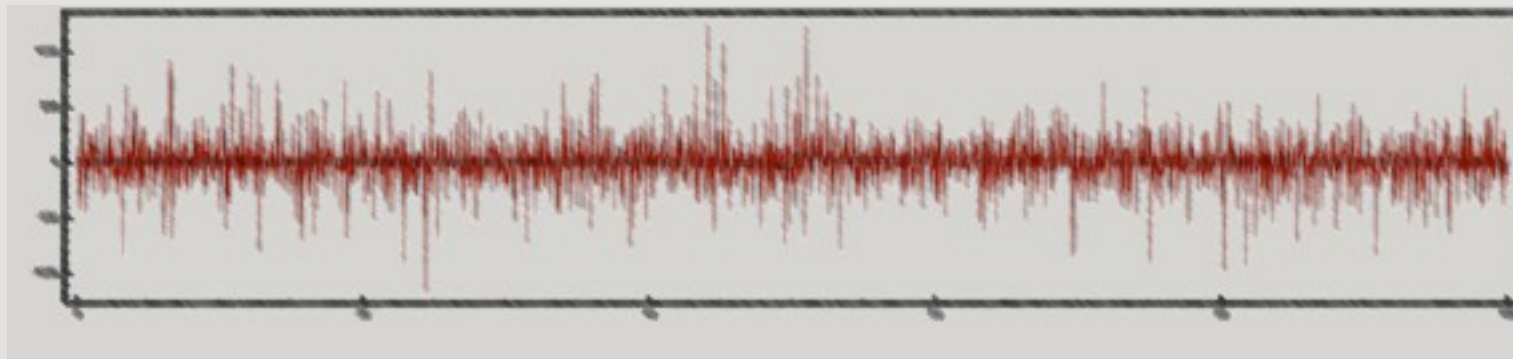
- ▷ **literature:** focus on fluctuation relations as rigorous bounds in nonequilibrium statistical mechanics
- ▷ **here:** FR as diagnostic tool in transport

transient fluctuation relations in transport

Chernogolovka, Sep. 20. 10

Boris Narozhny, Alessandro de Martino, Alexander Altland, Cologne University

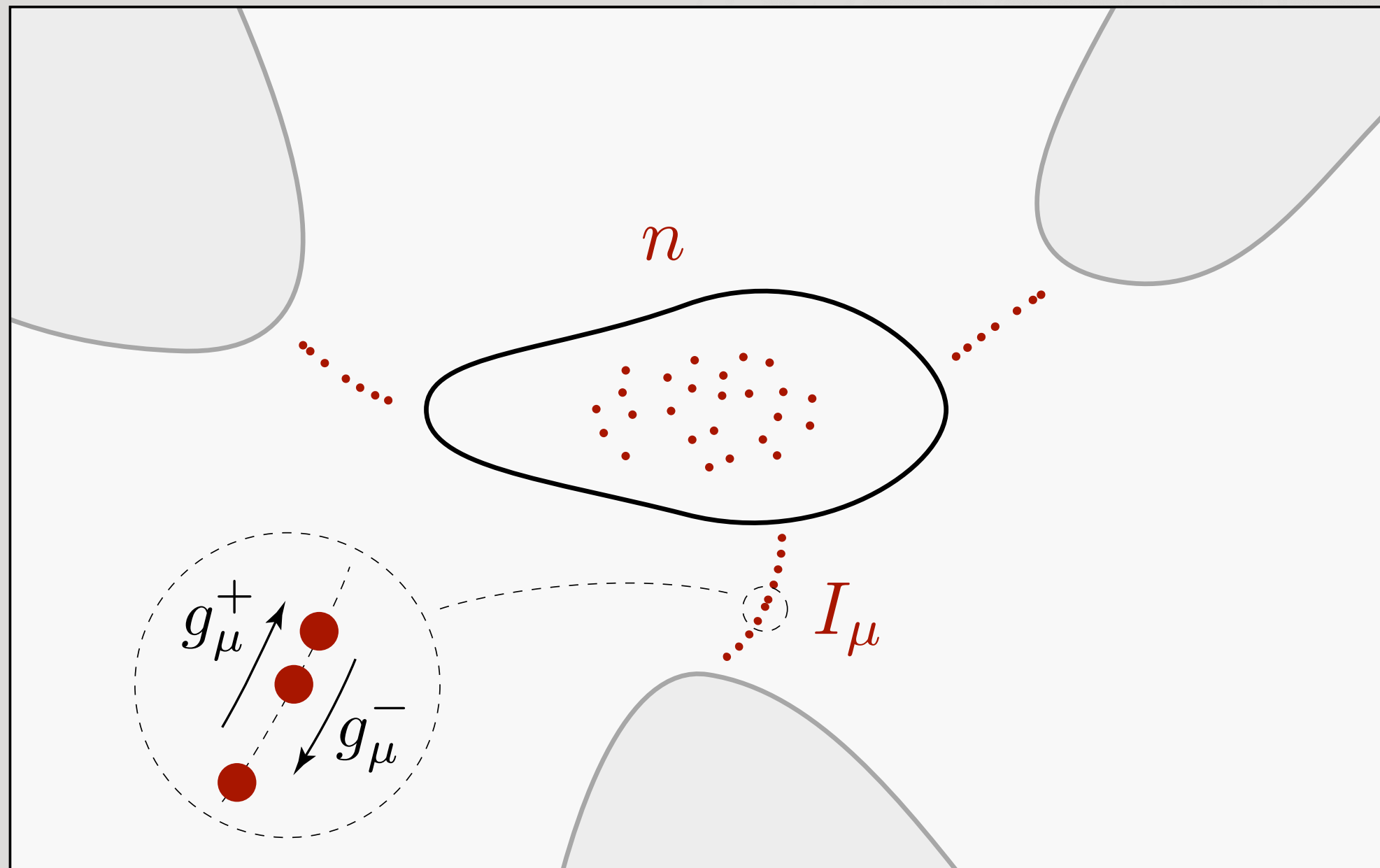
Reinhold Egger, Düsseldorf University



- ▷ master equation
- ▷ stochastic path integral
- ▷ fluctuation relations
- ▷ example: RC circuit

arXiv:1007.1826, 1005.4662

master equation and stochastic path integral



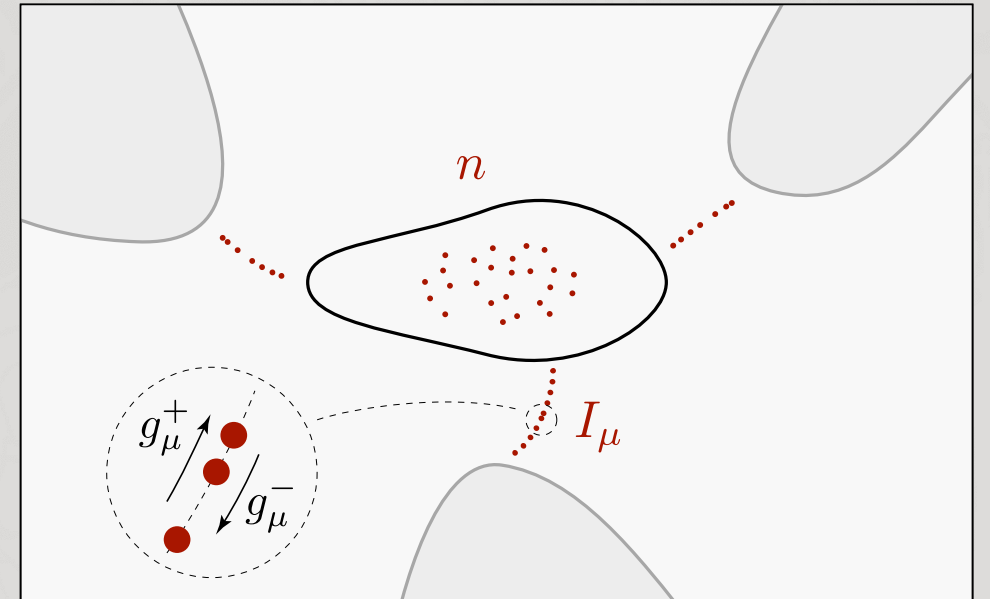
setup

▷ describe setup in terms of system **particle number** and **currents**

$$d_t n + \sum_{\nu=1}^M I_{\nu} = 0$$

quantity of interest:

$$\{I_1, \dots, I_{\mu}\} [g_1^{\pm}, \dots, g_{\mu}^{\pm}]$$



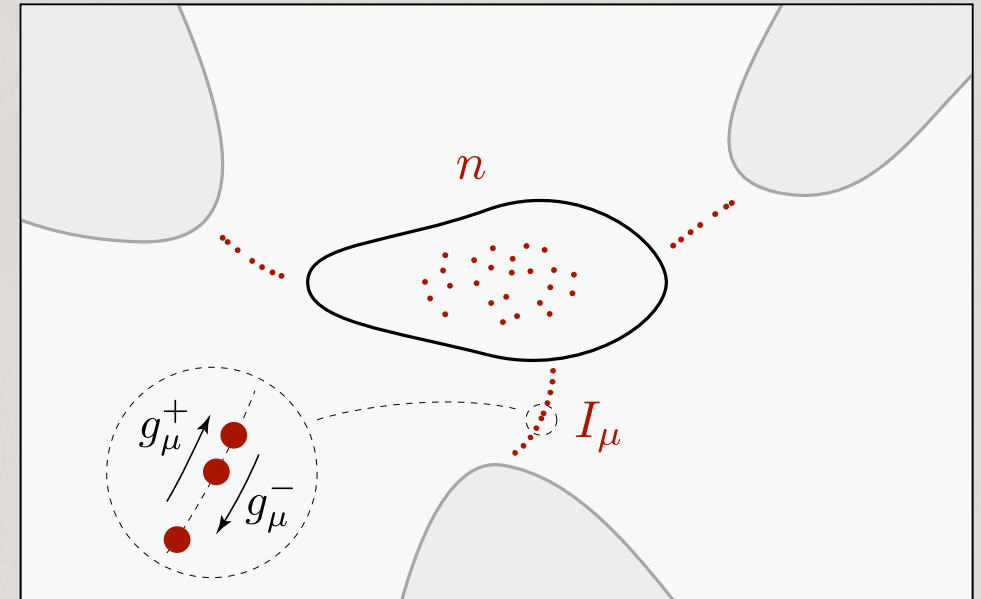
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- ▷ describe dynamics of particle currents in terms of **rates**

$$\frac{g_{\nu}^{+}}{g_{\nu}^{-}} = \exp(-\beta(\partial_n U - f_{\nu}))$$

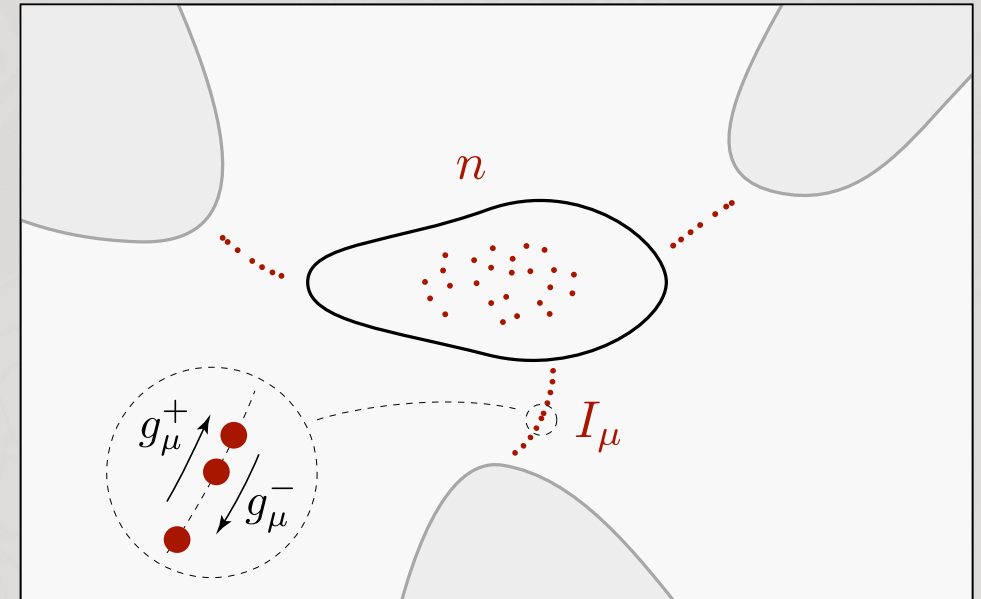
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- ▷ describe dynamics of particle currents in terms of **rates**

$$\frac{g_{\nu}^{+}}{g_{\nu}^{-}} = \exp(-\beta(\partial_n U - f_{\nu}))$$

- ▷ local energy change due to entering particle:

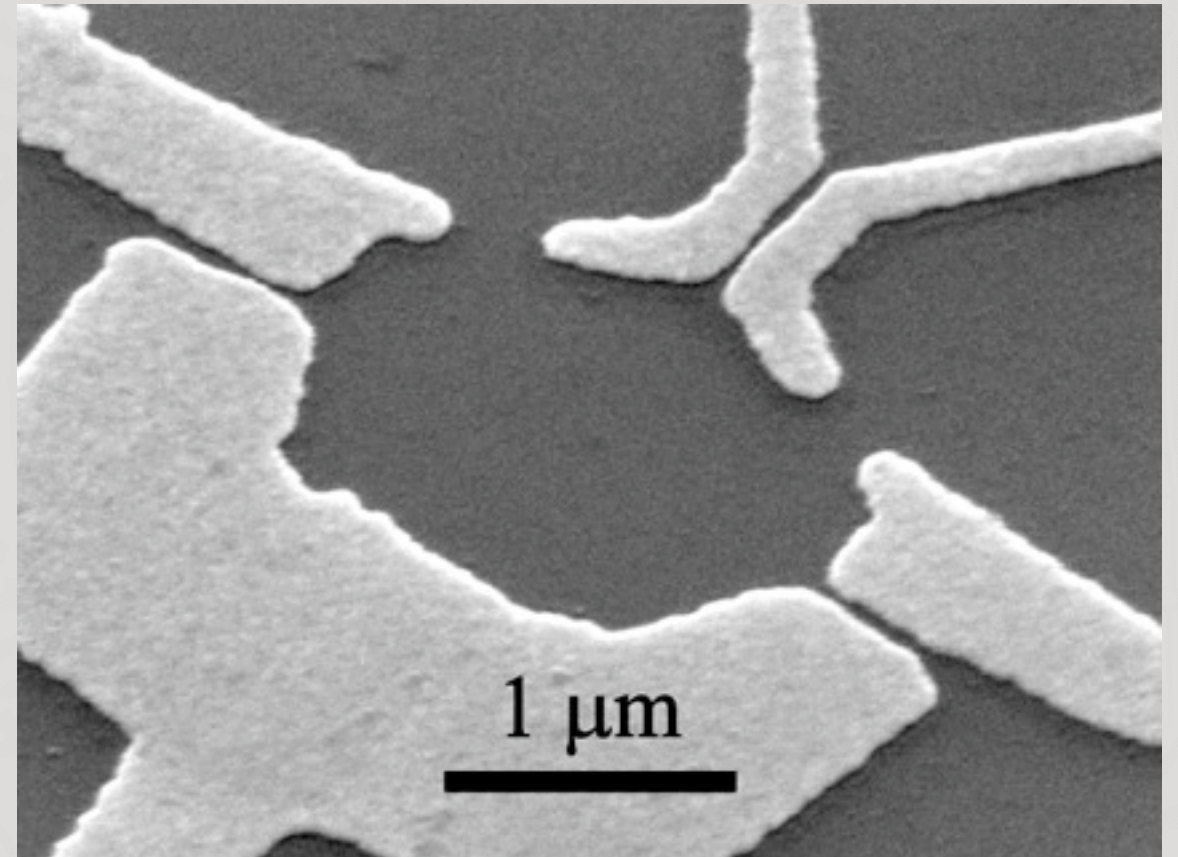
$$\Delta E_{\nu}(n) = U(n+1) - U(n) - f_{\mu} \simeq \partial_n U(n) - f_{\nu}$$

internal energy

driving force

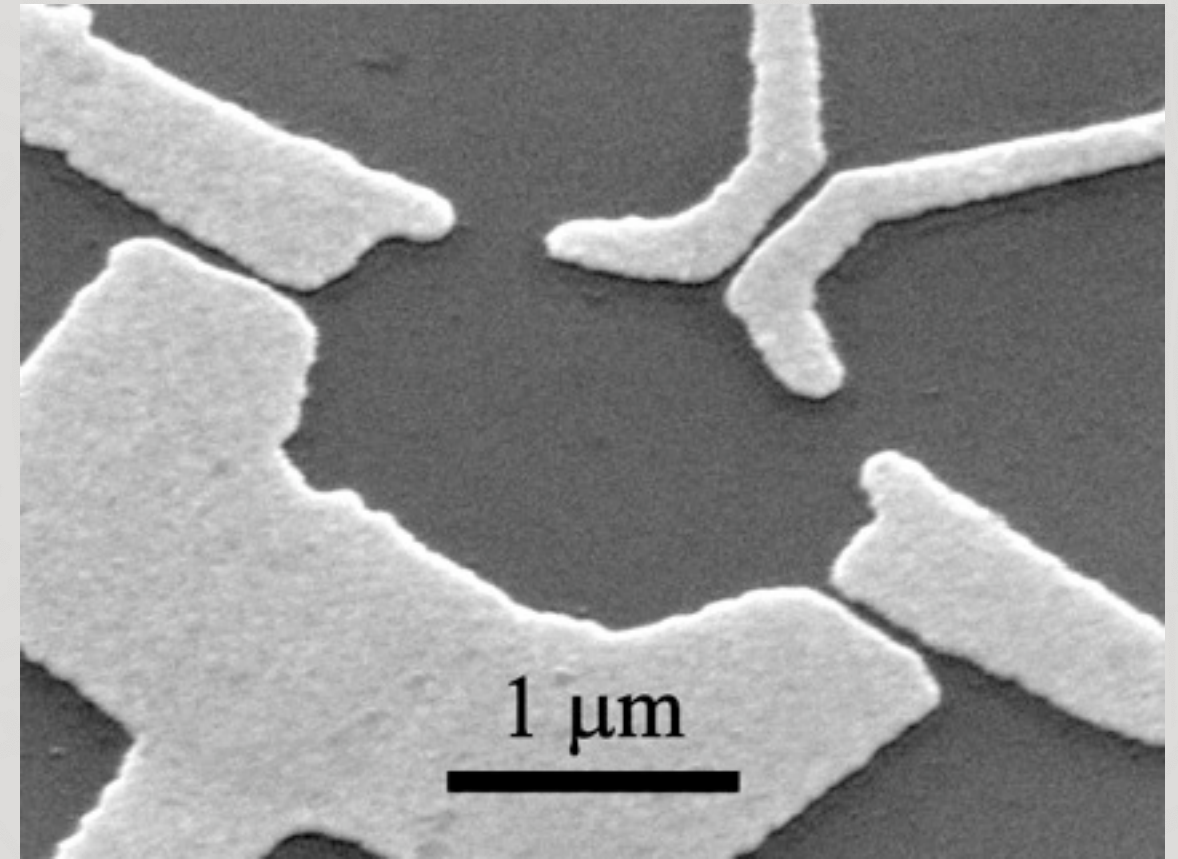
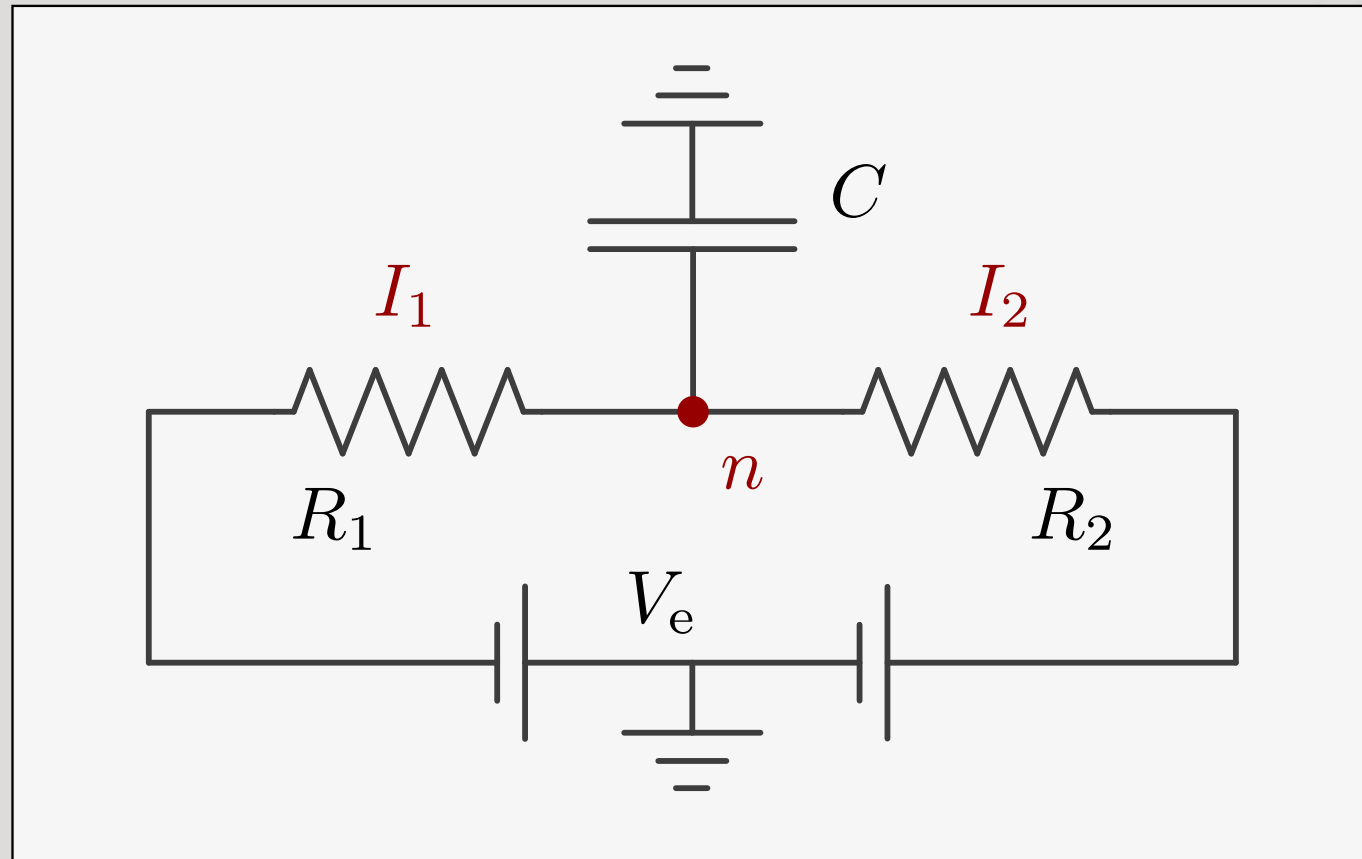
example

mesoscopic RC circuit



example

mesoscopic RC circuit



▷ island charge

▷ internal energy: $U(n) = \frac{n^2}{2C}$

▷ driving: $f_\nu = (-)^\nu \frac{V_e}{2}$

System	Variable n	$U(n)$	f_ν
electronic circuits (see Sec. IV)	charge	charging energy	bias voltages
molecular motors ³⁵	mechanochemical state of motor protein	load potential	ATP concentration
chemical reaction networks ³⁶	number of reaction partners	internal energy	chemostat concentrations
adaptive evolution ³⁷	allele frequencies	log equilibrium dist.	fitness gradients

master equation

▷ consider probability $P_t(n)$

$$\partial_t P_t(n) = -\hat{H}_g P_t(n),$$

$$\hat{H}_g(n, \hat{p}) = \sum_{\pm, \nu=1}^M (1 - e^{\mp \hat{p}}) g_{\nu, t}^{\pm}(n),$$

(one step) master equation

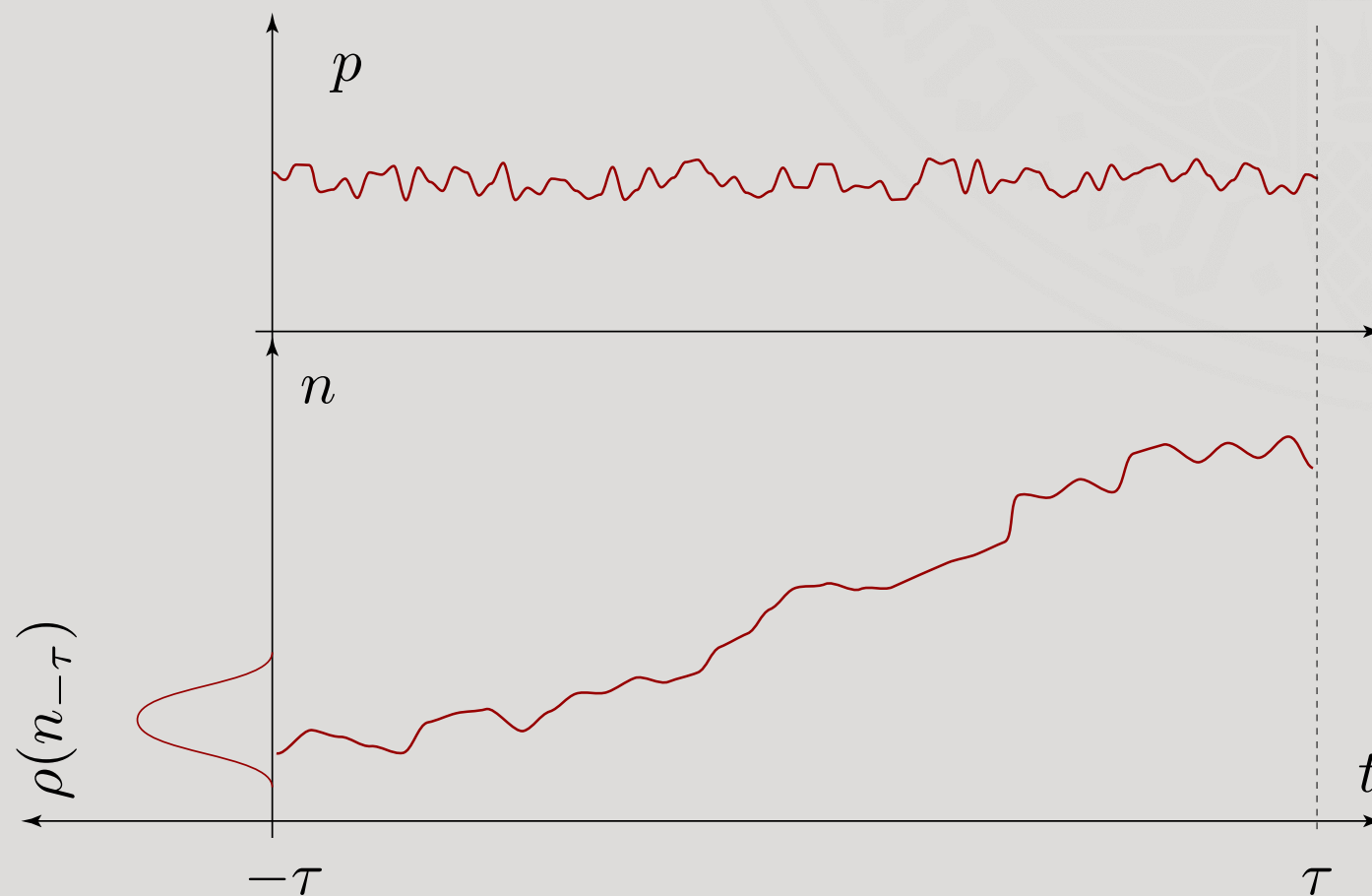
where $e^{\pm \hat{p}} f(n) = f((n \pm 1))$, $[\hat{p}, \hat{n}] = 1$

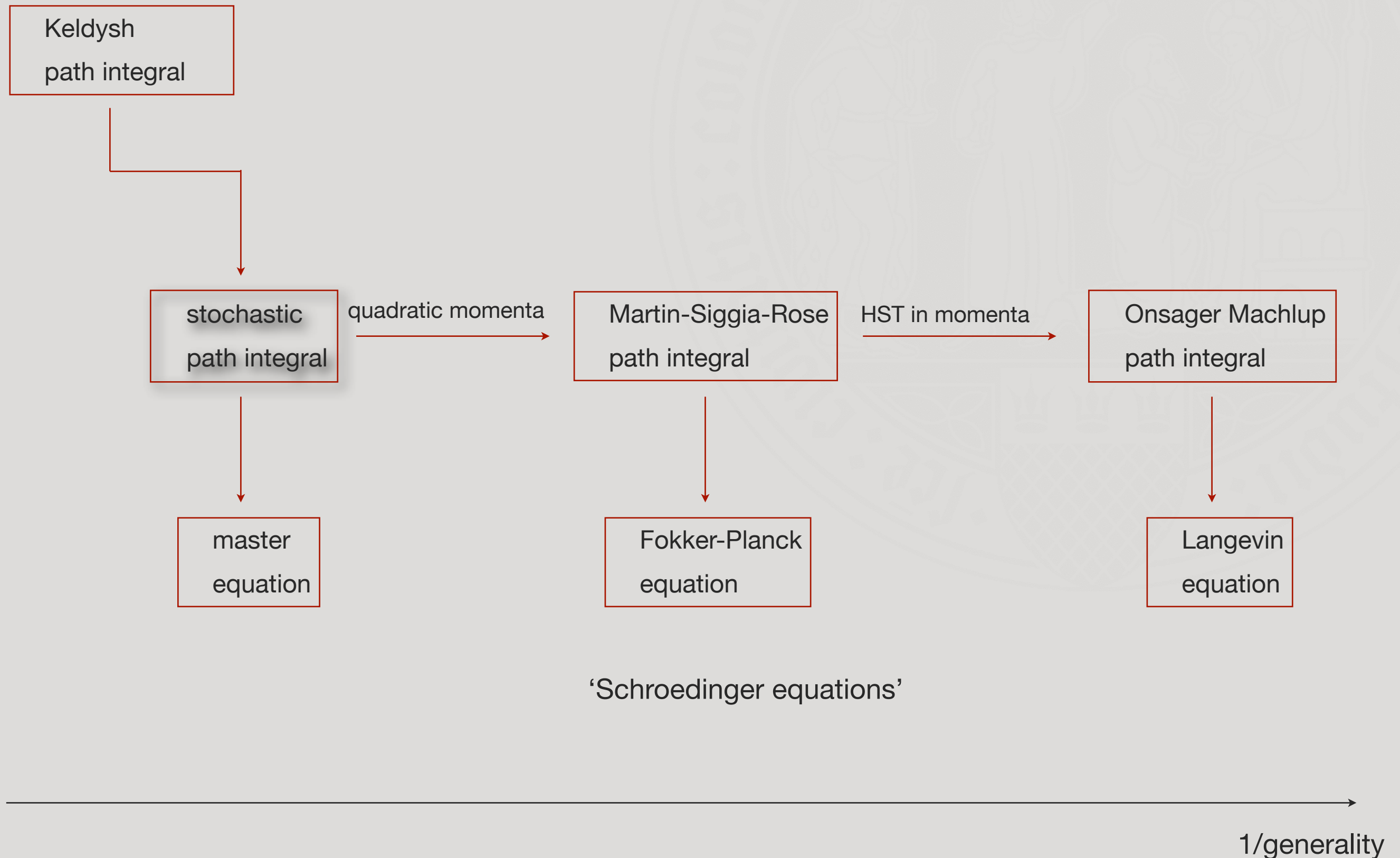
stochastic path integral

▷ represent probability in terms of (imaginary time) **stochastic path integral**

$$Z_f \equiv 1 = \sum_n P_t(n) = \int D(n, p) e^{-S_g[n, p]} \rho(n_{-\tau}),$$
$$S_g[n, p] = - \int_{-\tau}^{\tau} dt (p \dot{n} - H_g(n, p))$$

Kubo et al., 73





information from the path integral

▷ functional partition function normalized to unity. Want to obtain currents.

$$H_g(n, p) \rightarrow H_g(n, p, \chi) \equiv \sum_{\nu, \pm} \left(1 - e^{\mp(p - i\chi_\nu)} \right) g_\nu^\pm(n)$$

cf. vector potential

▷ generating functional (Pilgram et al. 03)

$$Z_f[\chi] = \int D(n, p) e^{-S_g[n, p, \chi]} \rho(n_{-\tau}),$$
$$S_g[n, p, \chi] \equiv - \int_{-\tau}^{\tau} dt (p\dot{n} - H_g(n, p, \chi)).$$

▷ current statistics

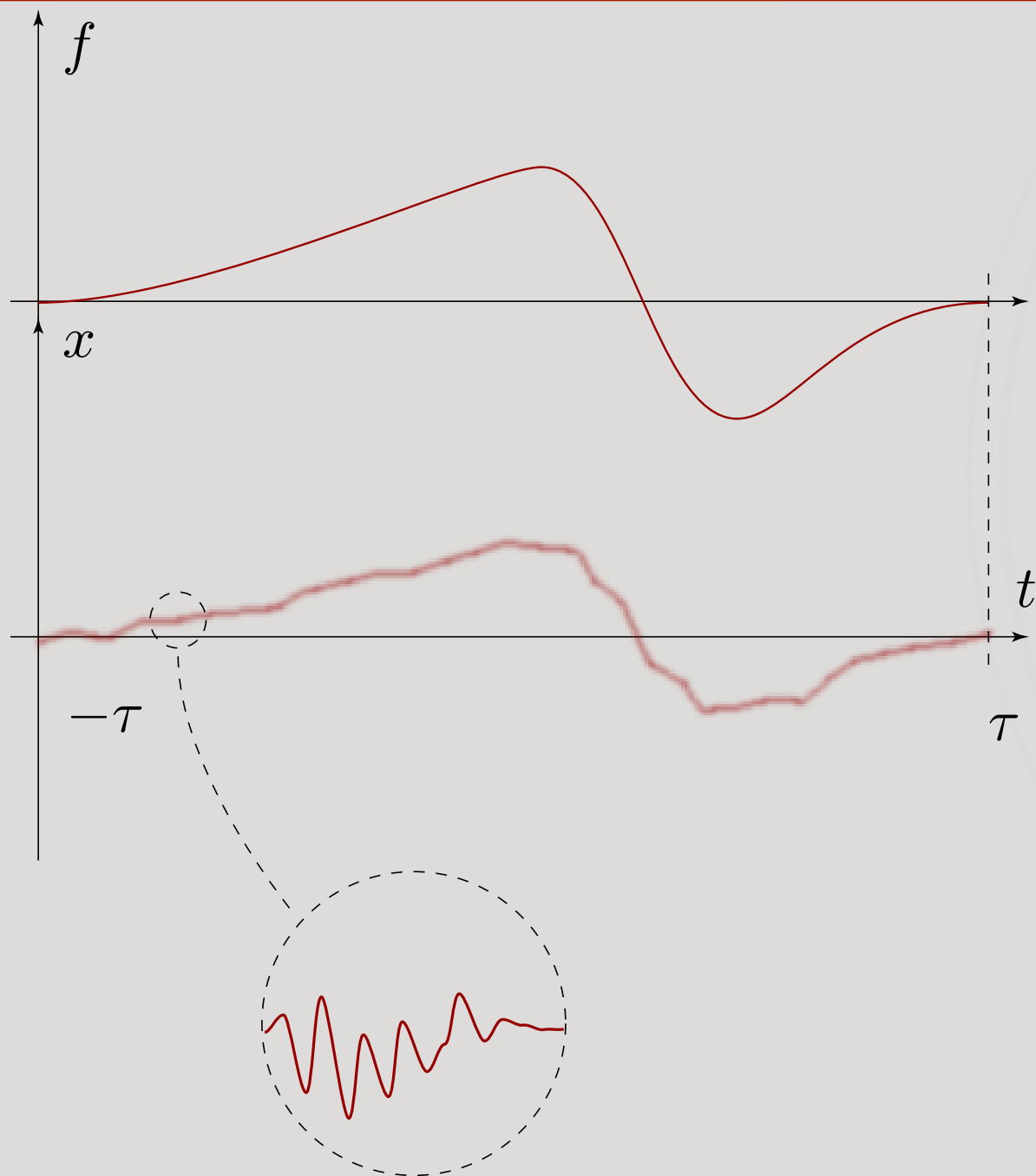
$$Z_f[\chi] = \left\langle e^{-i \sum_\nu \int_{-\tau}^{\tau} dt \chi_\nu I_\nu} \right\rangle_f,$$

$$\langle I_{\nu_1, t_1} I_{\nu_2, t_2} \dots \rangle = \frac{i\delta}{\delta \chi_{\nu_1, t_1}} \frac{i\delta}{\delta \chi_{\nu_2, t_2}} \dots \Big|_{\chi=0} Z_f[\chi] \Leftrightarrow$$

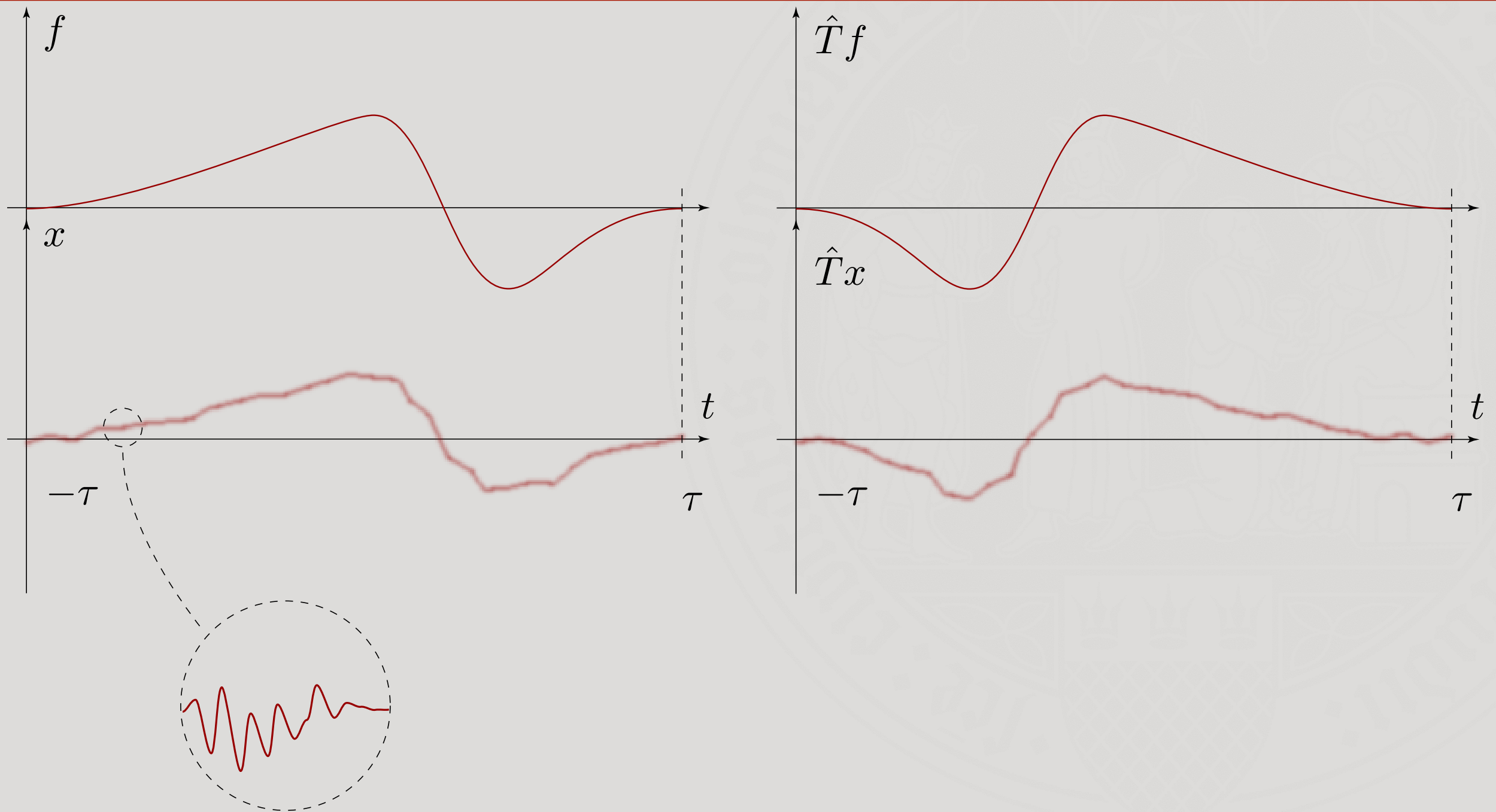


fluctuation relations

fluctuation relations

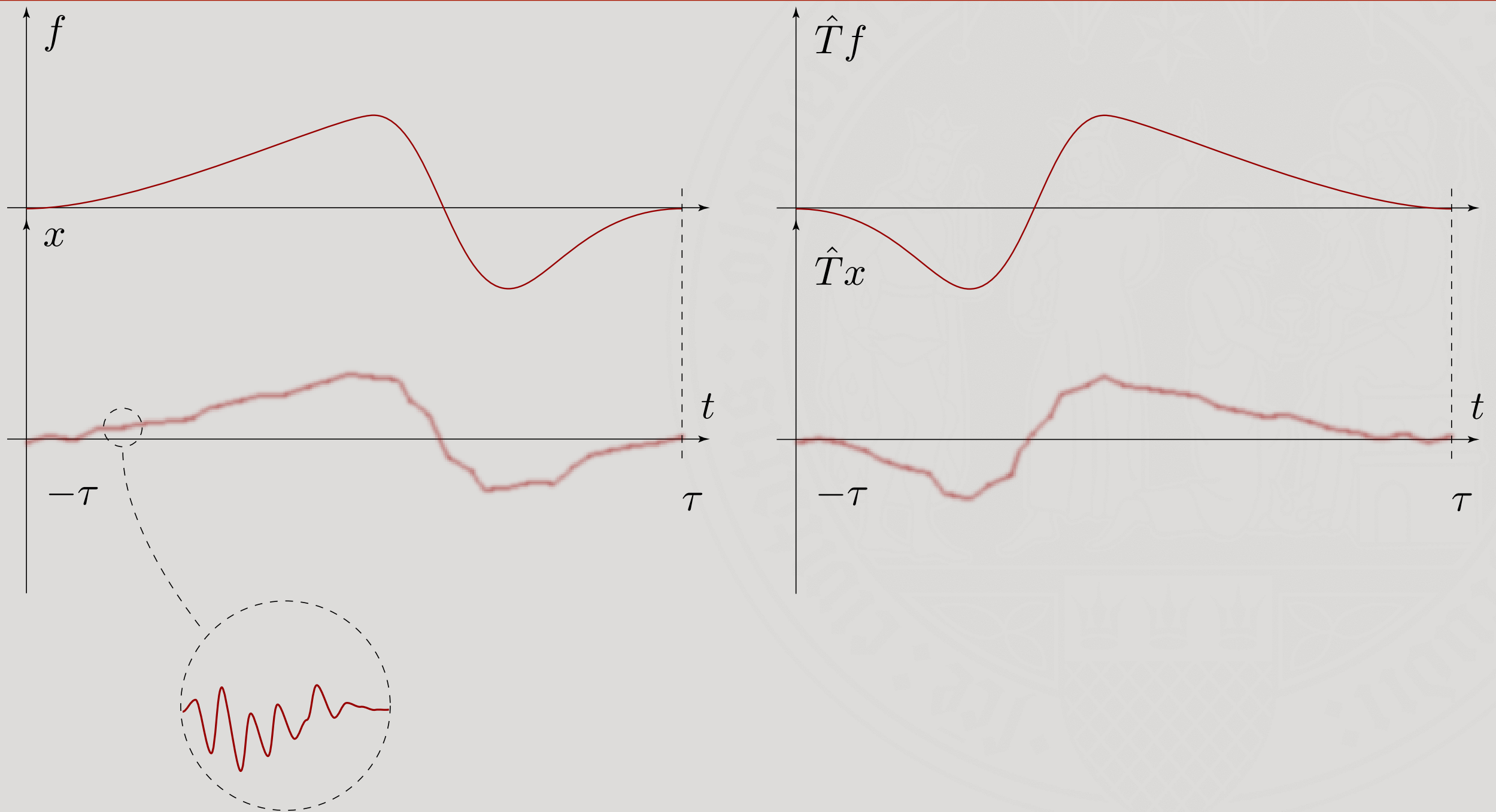


fluctuation relations



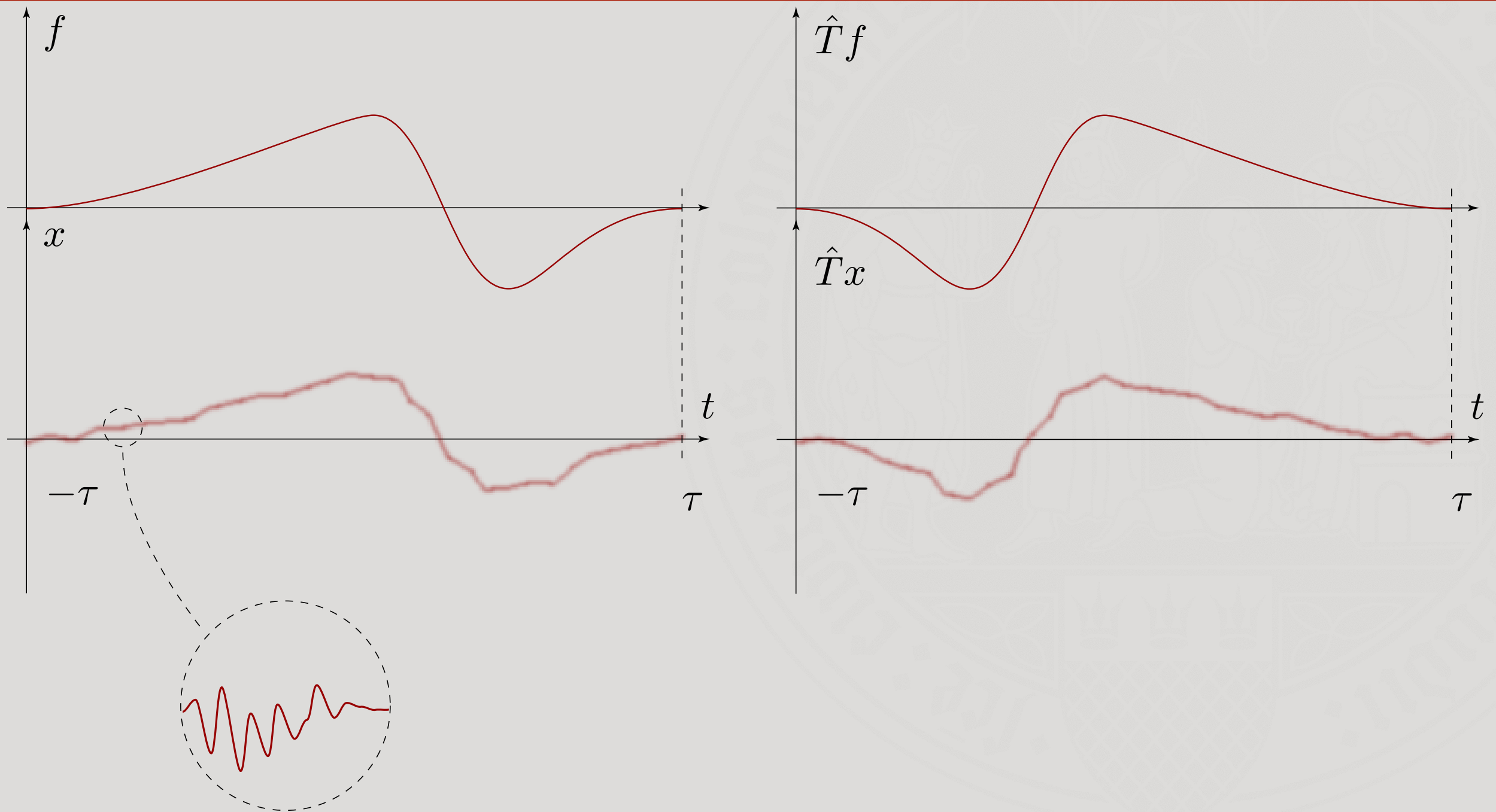
$$P_f[x] = P_b[\hat{T}x]$$

fluctuation relations



$$P_f[x] = P_b[\hat{T}x] \times e^{\beta Q[x]}$$

fluctuation relations



$$P_f[x] = P_b[\hat{T}x] \times e^{\beta Q[x]} \times e^{\beta(U[x] - \Delta F)} = P_b[\hat{T}x] \times e^{\beta(W[x] - \Delta F)}$$

$$S_g[n, p, \chi] = S_{\hat{T}g}[\hat{T}n, \hat{T}(p - \beta \partial_n U), \hat{T}(\chi + i\beta f)] + \beta[U(n_\tau) - U(n_{-\tau})].$$

$$X = n, g, U \quad X_t \rightarrow (\hat{T}X)_t = X_{-t}$$

$$X = p, \chi_\nu, I_\nu \quad X_t \rightarrow (\hat{T}X)_t = -X_{-t}$$

$$S_g[n, p, \chi] = S_{\hat{T}g}[\hat{T}n, \hat{T}(p - \beta \partial_n U), \hat{T}(\chi + i\beta f)] + \beta[U(n_\tau) - U(n_{-\tau})].$$

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$$Z_f[\chi] = \int D(n, p) e^{-S_{\tau_g}[n, p, \mathcal{T}(\chi + i\beta f)]} \rho(n_{-\tau}) = Z_b[\mathcal{T}(\chi + i\beta f)].$$

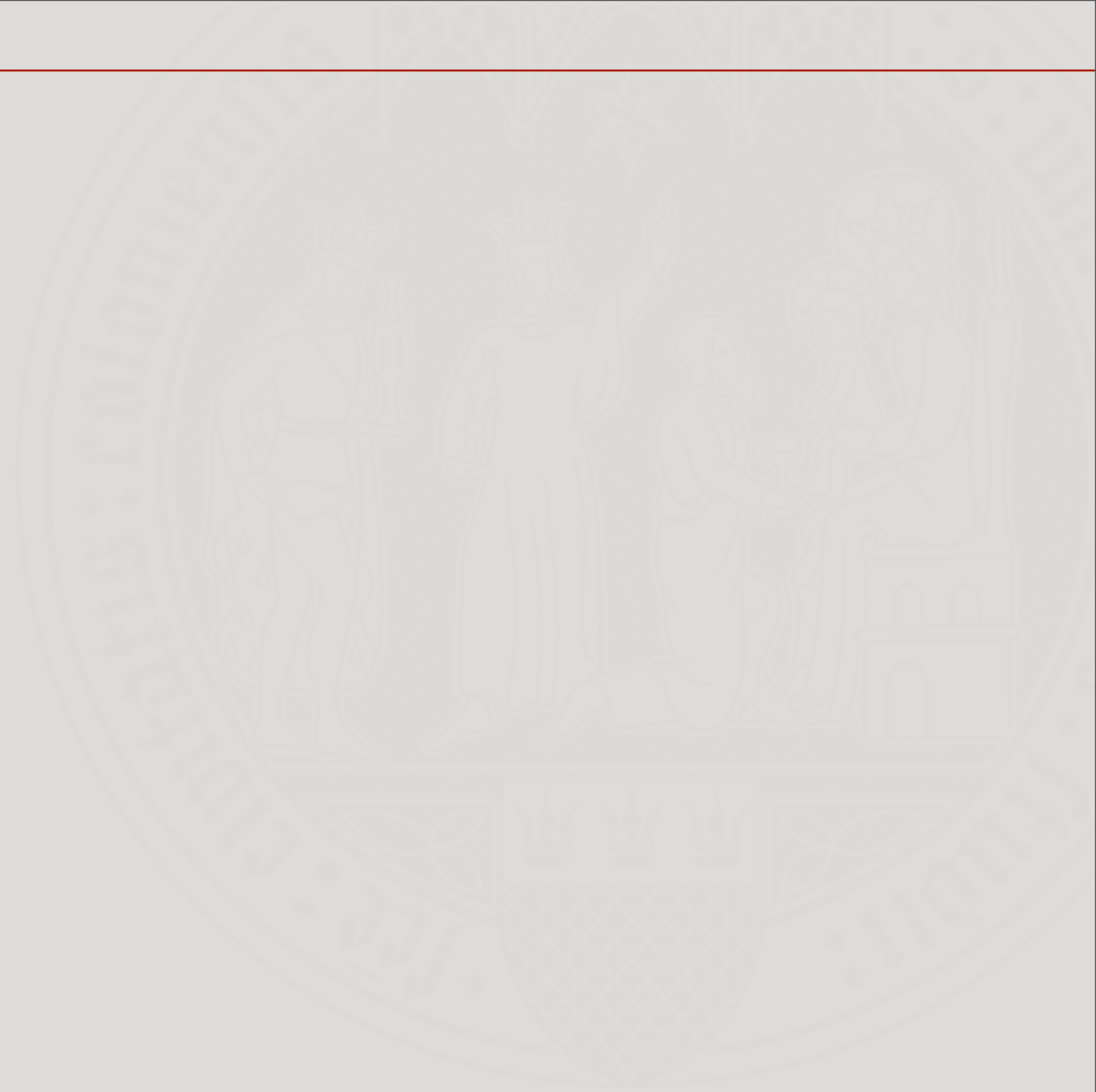
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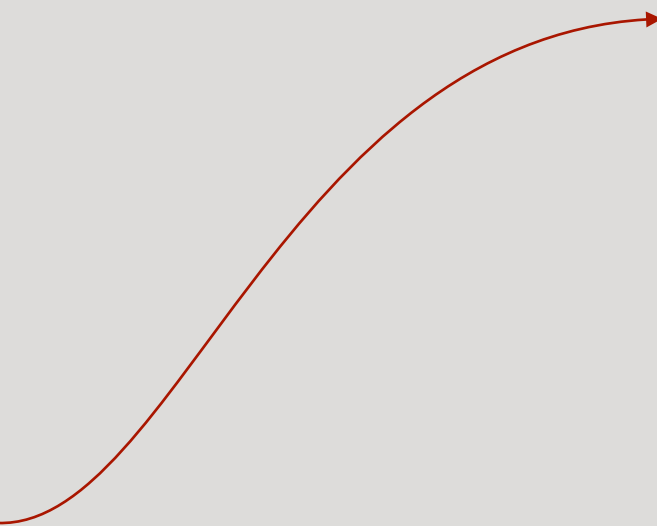
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fluctuation relations




$$P_f[I] = P_b[\hat{T}I] e^{-\beta \int_{-\tau}^{\tau} dt \sum_{\nu} f_{\nu} I_{\nu}}$$

functional Crooks relation (cf. Bochkov & Kusovlev, 81)

$$\left\langle e^{\beta \int_{-\tau}^{\tau} dt \sum_{\nu} f_{\nu} I_{\nu}} \right\rangle_f = 1$$

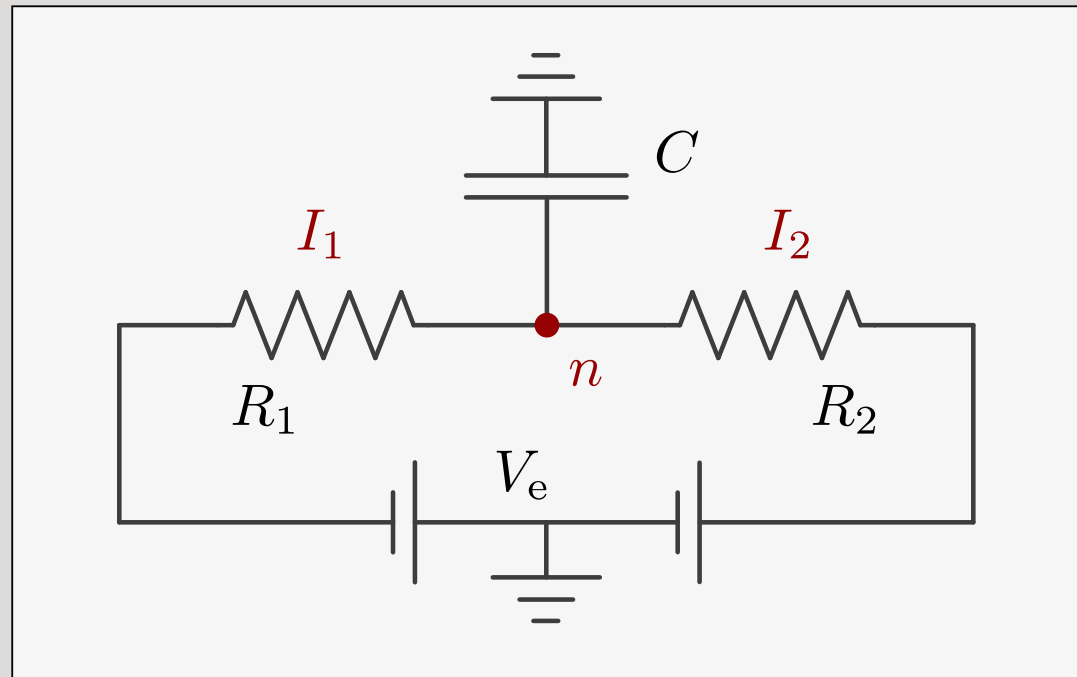
functional Jarzynski relation

applications

- ▷ information on rare event statistics
- ▷ generalized Onsager relations for AC transport
- ▷ information on effective distributions prevalent in the system
- ▷ information on joint statistics of particle transport and energy

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application: mesoscopic RC circuit



▷ two terminal setup

▷ system parameters: $G_\nu = 20 \frac{e^2}{h}$

V fluctuates at scales larger than RC-time

V/T large (shot noise regime) or small (thermal regime)

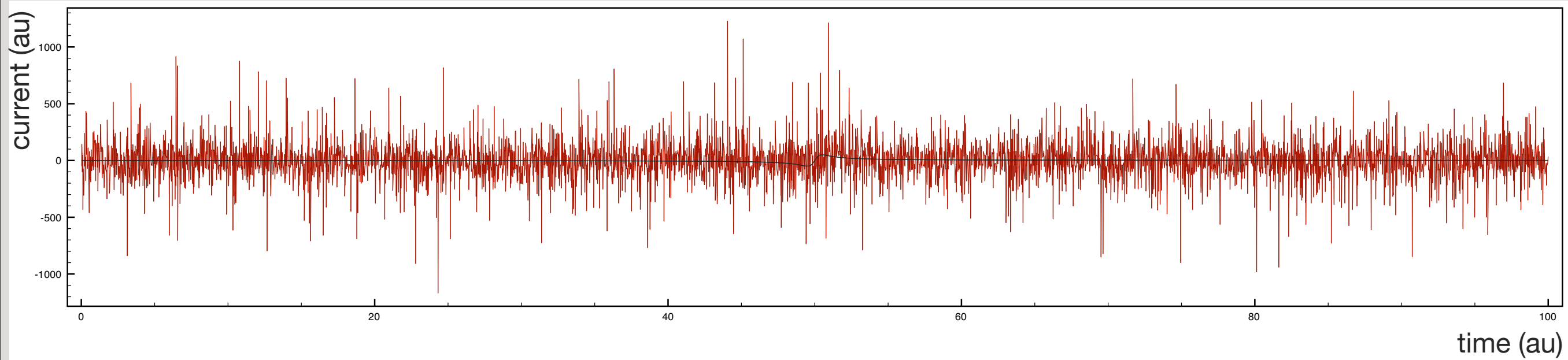
▷ microscopic description through master equation with rates

$$g_{\nu,t}^{\pm}(n) = G_\nu \frac{\pm \kappa_{\nu,t}(n)}{e^{\pm \beta \kappa_{\nu,t}(n)} - 1}, \quad \kappa_{\nu,t}(n) = \partial_n U(n) + (-)^\nu \frac{V_t}{2}.$$

applications I: rare event statistics

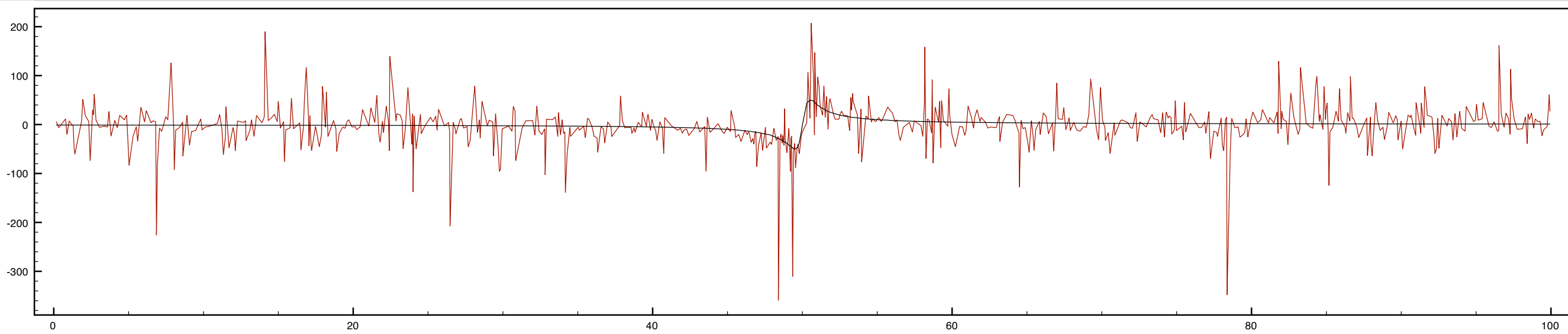
current fluctuations

▷ near equilibrium $T/V \gg 1$: thermal (Johnson-Nyquist) fluctuations



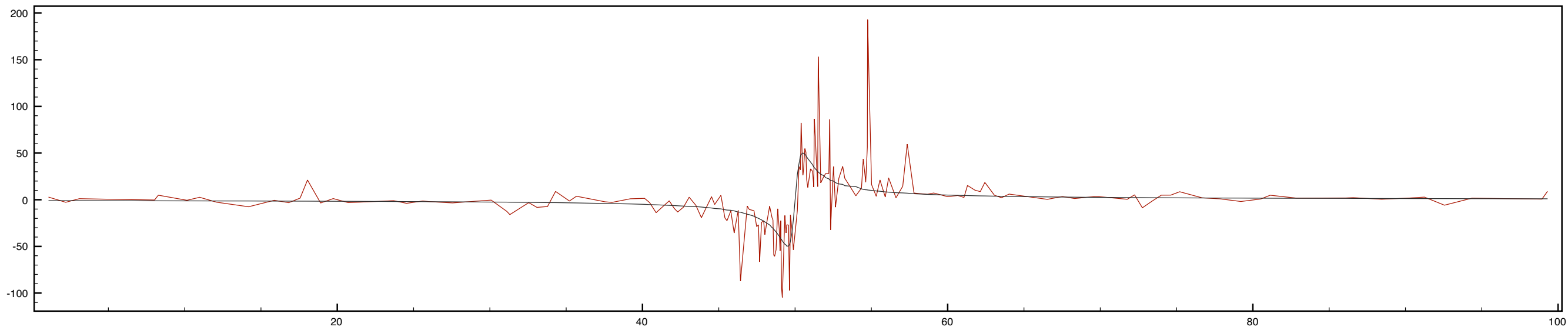
current fluctuations

▷ crossover regime $T/V \simeq 1$



current fluctuations

▷ out of equilibrium $T/V \ll 1$



current fluctuations vs. Jarzynski relation

$$\left\langle e^{\beta \int_{-\tau}^{\tau} dt V(I_2 - I_1)} \right\rangle_f = 1$$

functional Jarzynski relation

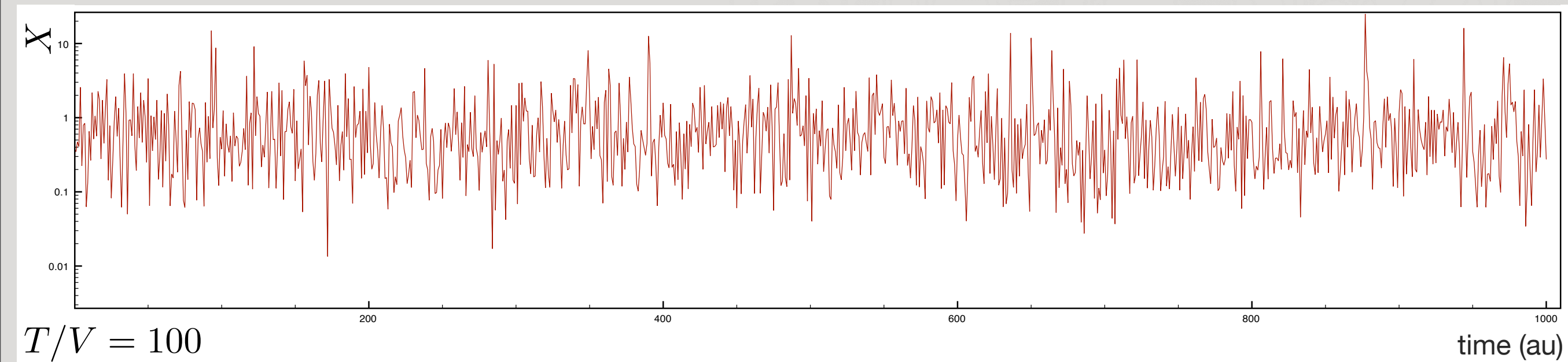
off equilibrium: probes rare events through variable $X \equiv e^{\beta \int_{-\tau}^{\tau} dt V(I_2 - I_1)}$

current fluctuations vs. Jarzynski relation

$$\left\langle e^{\beta \int_{-\tau}^{\tau} dt V(I_2 - I_1)} \right\rangle_f = 1$$

functional Jarzynski relation

off equilibrium: probes rare events through variable $X \equiv e^{\beta \int_{-\tau}^{\tau} dt V(I_2 - I_1)}$

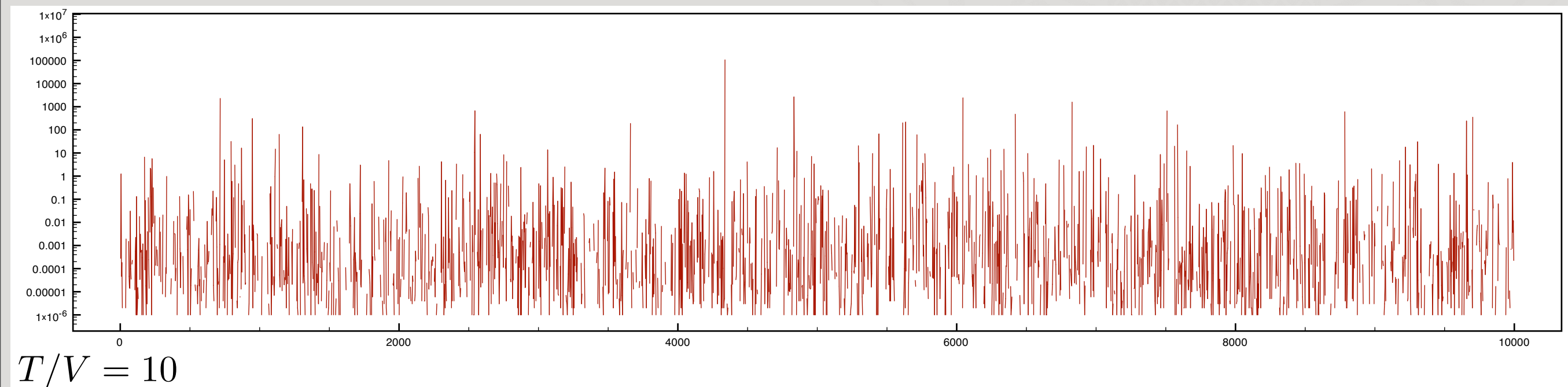
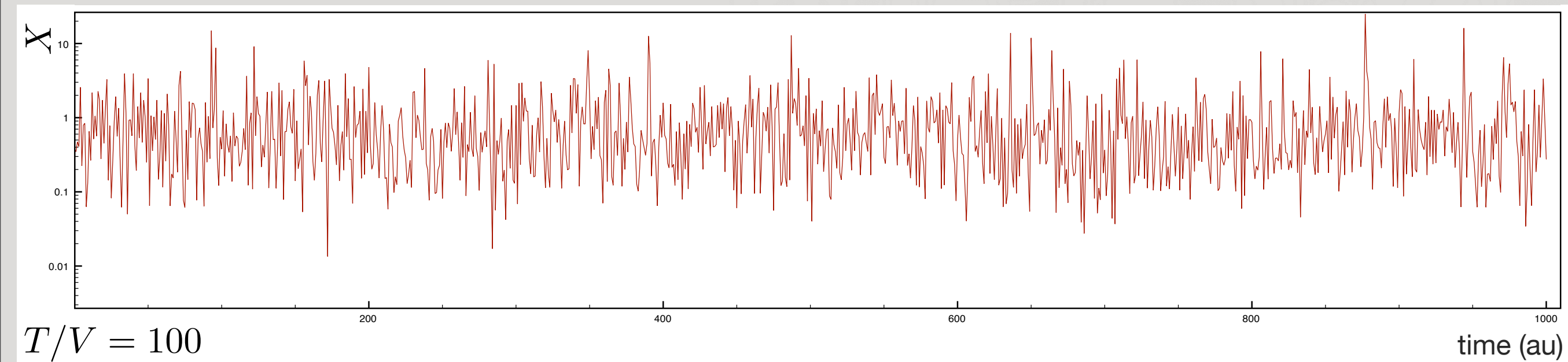


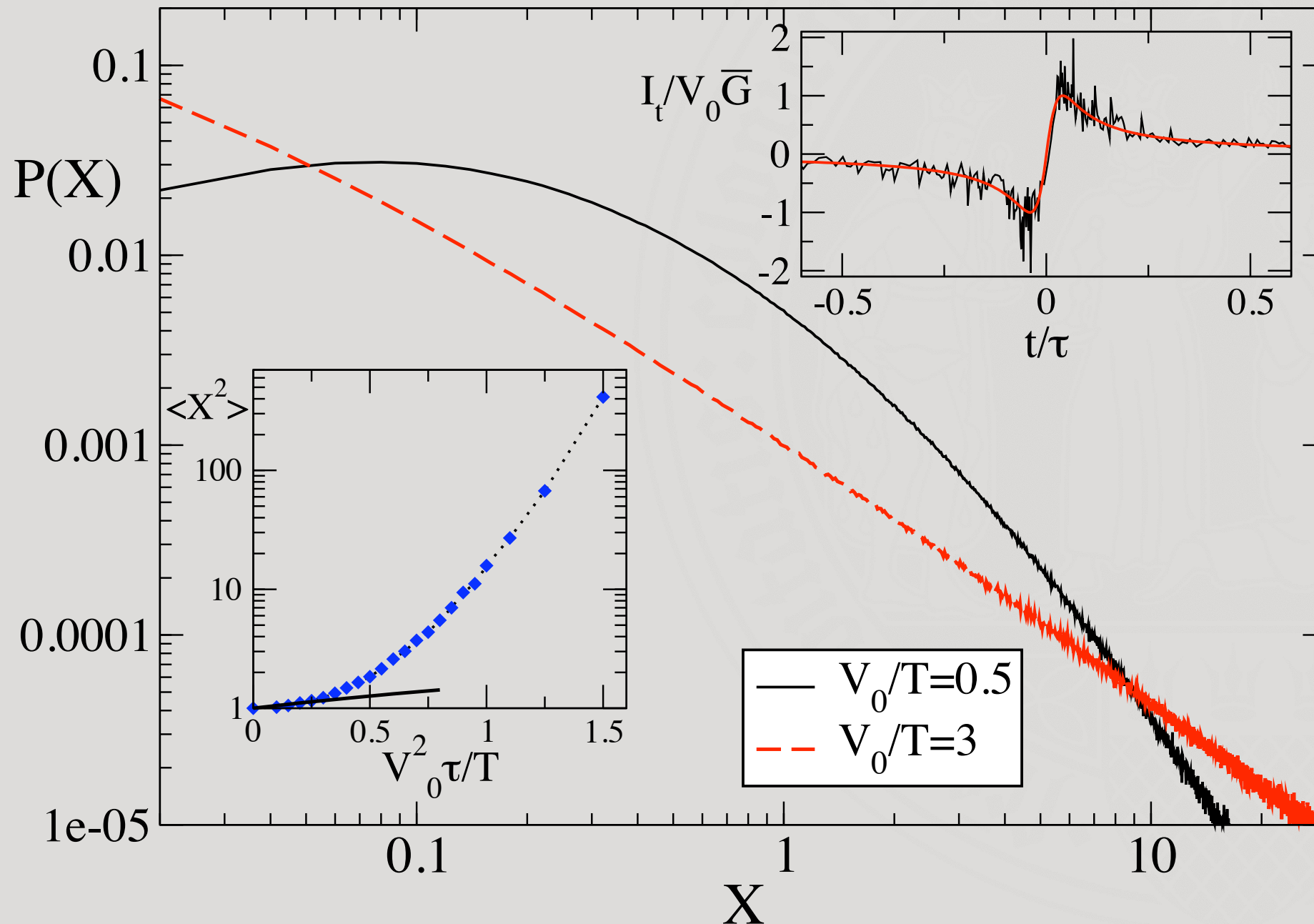
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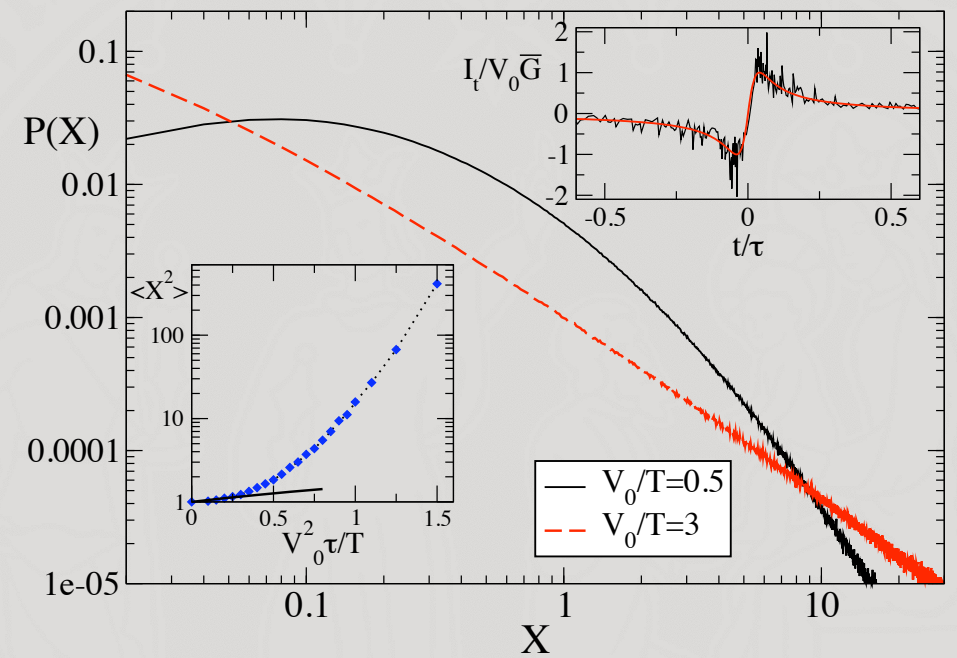




path integral can be applied to estimate fluctuations of X

thermal regime: $\langle X^2 \rangle_f \sim e^{\tau \beta (f_1 - f_2)^2 \frac{g_1 g_2}{g_1 + g_2}}$

shot noise regime: $\langle X^2 \rangle_f \sim e^{4\tau |\langle I_1 \rangle \langle I_2 \rangle|^{1/2} e^{\beta |f_1 - f_2|/2}}$



optimal fluctuation approach to path integral

thermal: Gaussian fluctuations around mean

$$\langle X^2 \rangle_f \sim e^{\tau \frac{V^2}{T} \frac{g_1 g_2}{g_1 + g_2}} = e^{\tau \langle I \rangle \frac{V}{T}}$$

shot noise: rare event statistics

$$\langle X^2 \rangle_f \sim e^{4\tau |\langle I_1 \rangle \langle I_2 \rangle|^{\frac{1}{2}}} e^{\frac{V}{2T}}$$

fluctuation statistics of X sensitive probe of nonequilibrium transport processes

applications II: distribution diagnostics

use fluctuation relations to explore state of the system

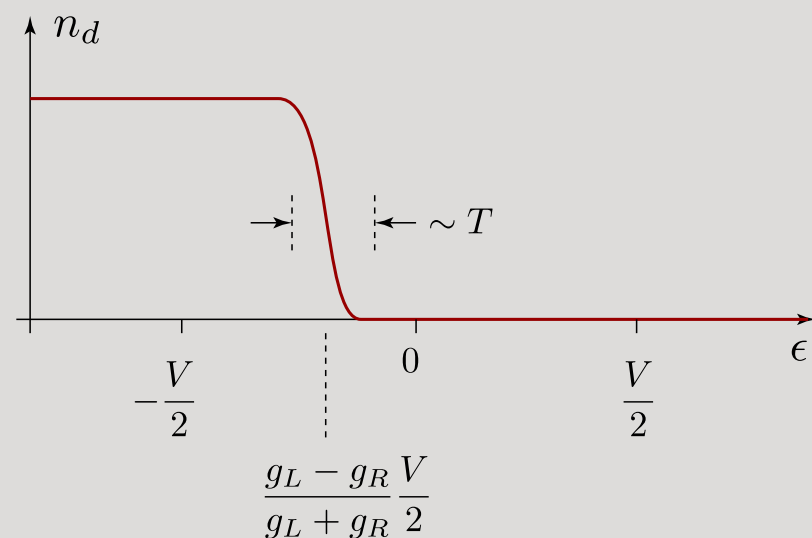
derivation of FRs based on balance relation $\frac{g_{\nu}^{+}}{g_{\nu}^{-}} = \exp(-\beta(\partial_n U - f_{\nu}))$

model dot distribution functions

use fluctuation relations to explore state of the system

derivation of FRs based on balance relation $\frac{g_{\nu}^{+}}{g_{\nu}^{-}} = \exp(-\beta(\partial_n U - f_{\nu}))$

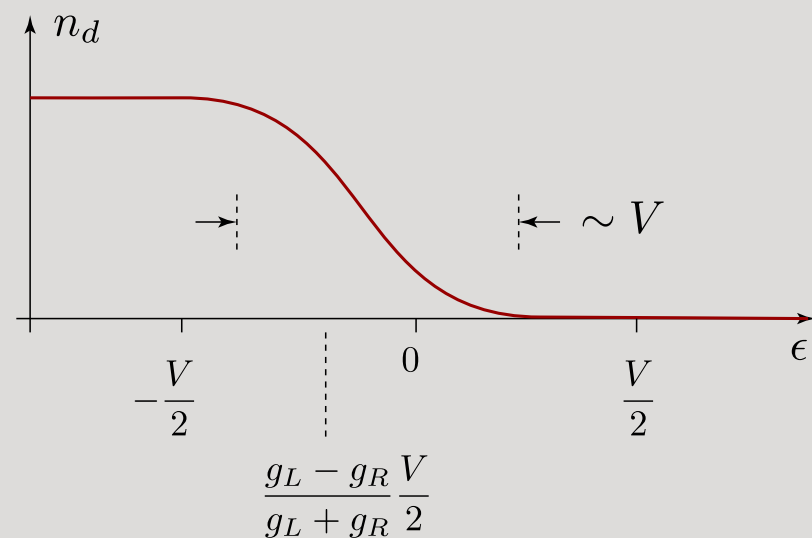
model dot distribution functions



$\tau_{ee} < \tau_d$, cooled

+

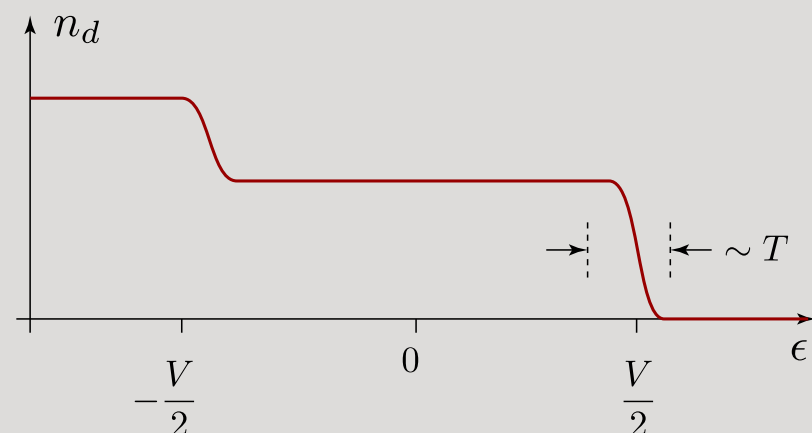
‘standard’ FR



$\tau_{ee} < \tau_d$, isolated

-

effective or extended FR



$\tau_{ee} > \tau_d$

(+)

standard FR

Effective fluctuation relations

▷ breakdown of $\frac{g_\nu^+}{g_\nu^-} = \exp\left(\frac{1}{T}(\partial_n U - f_\nu)\right)$



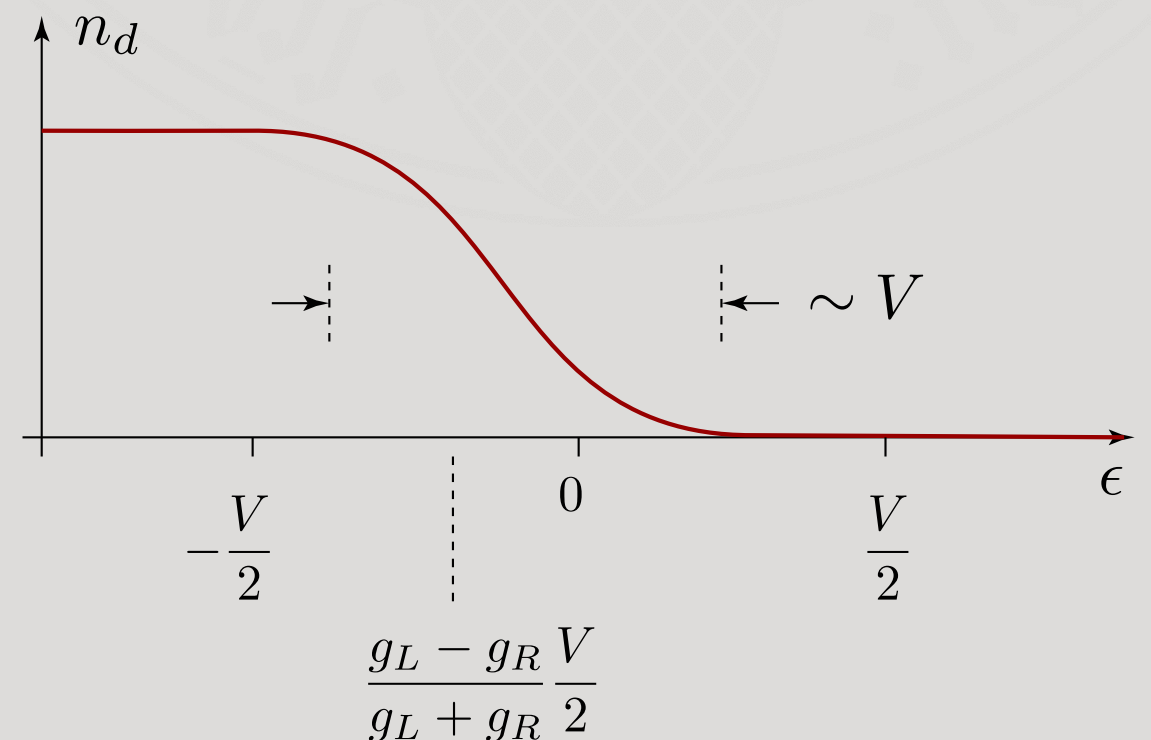
▷ need to describe transport at higher resolution (e.g. particle & energy transport)

▷ *define* effective temperature through

$$\frac{g_\nu^+}{g_\nu^-} = \exp\left(\frac{1}{T_{\text{eff},\nu}}(\partial_n U - f_\nu)\right)$$

phenomenology

comparison to (measured) fluctuation statistics contains information on mechanisms of nonequilibrium driving.



fluctuation relations in stochastic dynamics

- ▷ are useful when included in **out of equilibrium theory of transport**
- ▷ enable one to describe **statistics of rare events**

